

"Self-concordant Schrödinger operators"

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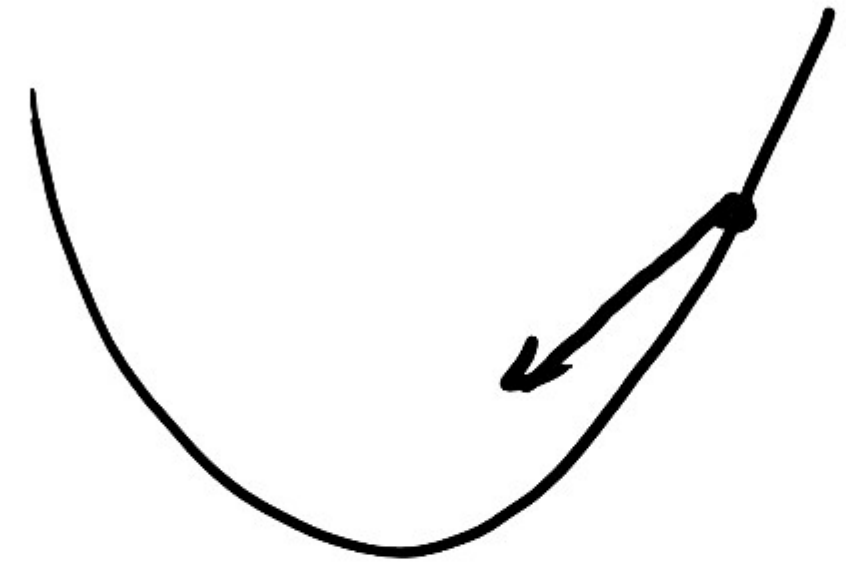
(CNRS, Université Paris Cité, IRIF)

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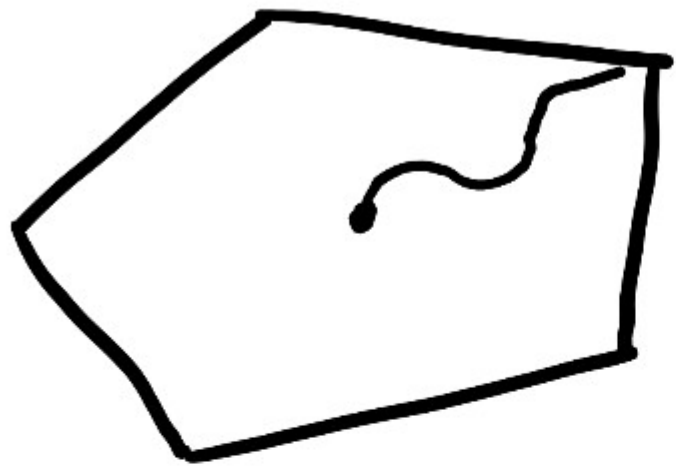
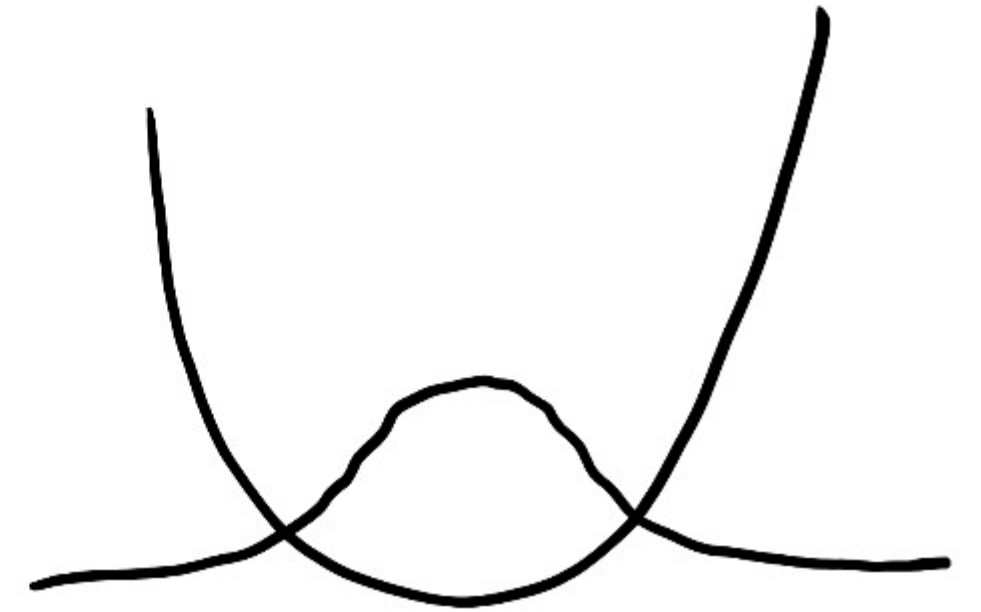
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OUTLINE

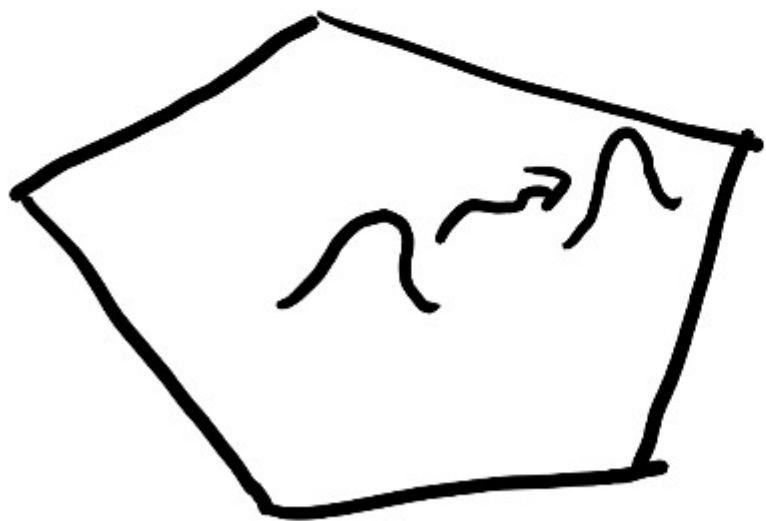
- (2nd order) optimization



- (2nd order) quantum optimization



- interior point methods



- quantum interior point method

(1st order) optimization

convex, smooth $f: \mathbb{R}^n \rightarrow \mathbb{R}$

"1st order" information
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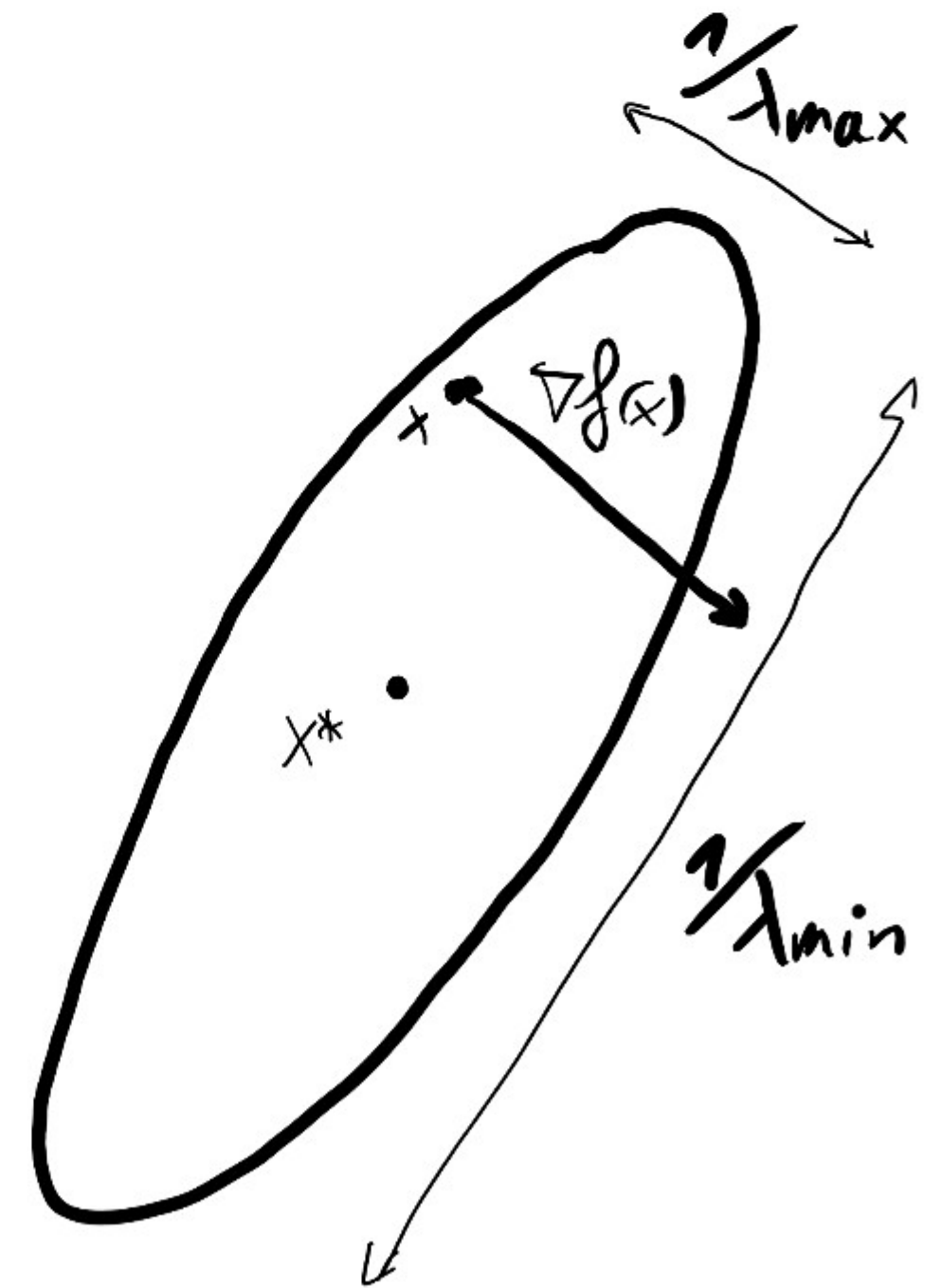
$\leadsto$ gradient descent:

$$x' = x - \eta \nabla f(x)$$

base case:

quadratic $f(x) = \frac{1}{2} x^T A x$, $A > 0$

$$! \text{ \# steps } \sim \frac{\lambda_{\max}(A)}{\lambda_{\min}(A)} = \kappa(A)$$



2nd order optimization:

Hessian $H(x) = \Delta^2 f(x) > 0$

smooth, (strongly) convex $f: \mathbb{R}^n \rightarrow \mathbb{R}$

↳ Newton step:

$$x' = x - \eta H(x)^{-1} \nabla f(x)$$

↳ "2nd order" information

e.g., quadratic $f(x) = \frac{1}{2} x^T A x$:

$$x' = x - \eta A^{-1} (Ax) = x - \eta x \quad \leadsto \# \text{ steps} \sim 1$$

2nd order optimization:

$$x' = x - \eta H(x)^{-1} \nabla f(x)$$

? general convex f

— f is **SELF-CONCORDANT** if $H(x)$ is "stable":

$$\| \delta \|_x \leq \varepsilon \implies H(x + \delta) \approx_\varepsilon H(x)$$

$(\| \delta \|_x = \sqrt{\langle \delta, H(x) \delta \rangle})$

— [Nesterov-Nemirovski: '94]: f self-concordant and x suff. close

$\left. \begin{array}{l} f \text{ self-concordant,} \\ x_0 \text{ "sufficiently close" to } x^* \end{array} \right\} \implies f(x_k) - f(x^*) \leq \frac{1}{2^k}$ — ! independent of
Cond. number
(and dimension)

"(1st order) Quantum optimization"

Schrödinger operator

$$\mathcal{H} = -\frac{1}{2}\Delta + f$$

smooth, convex
on $\text{dom}(\mathcal{H}) \ni C_c^\infty(\mathbb{R}^n)$

— if $f \geq 0$ and $f(x) \rightarrow \infty$ as $x \rightarrow \infty$:

ground state energy

$\mathcal{H}$ has discrete, nonnegative spectrum $0 \leq \lambda_0 \leq \lambda_1 \leq \lambda_2 \leq \dots$

$\lambda_1 - \lambda_0 = \text{spectral gap}$
 $\delta(\mathcal{H})$

— key example: $\mathcal{H}_0 = -\frac{1}{2}\Delta + \frac{1}{2}x^T A x$

= quantum harmonic oscillator



g.s. energy

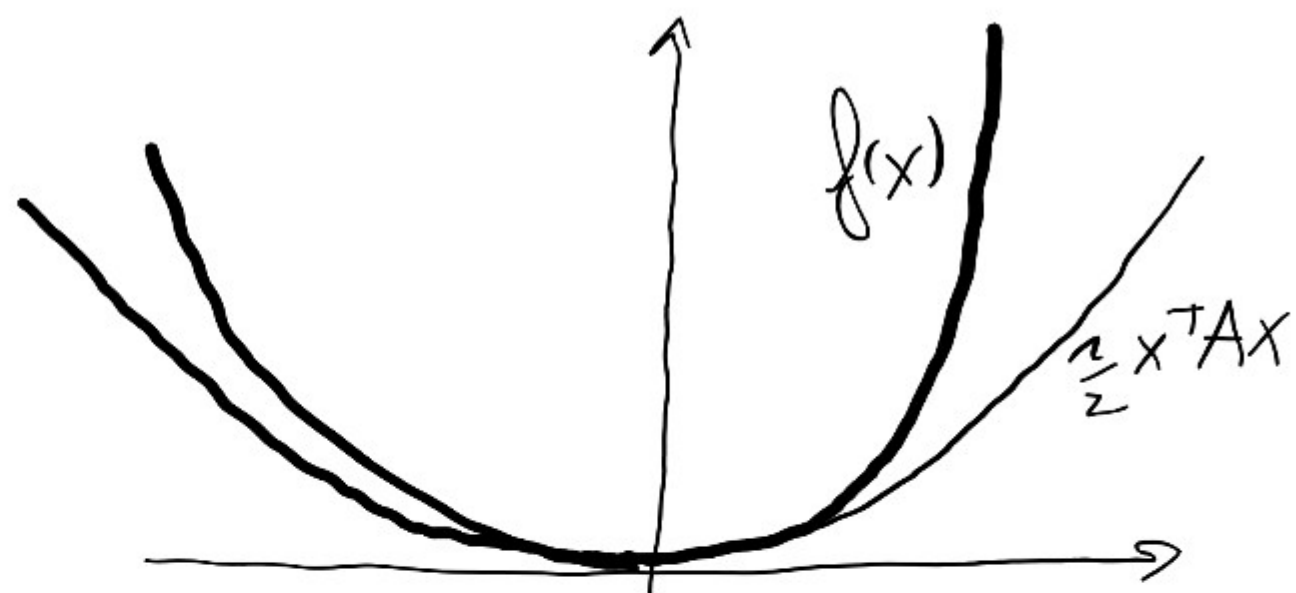
$$\lambda_0(\mathcal{H}_0) = \frac{1}{2} \text{Tr}[\sqrt{A}]$$

$$\lambda_1(\mathcal{H}_0) = \frac{1}{2} \text{Tr}[\sqrt{A}] + \lambda_{\min}[\sqrt{A}]$$

$$\frac{\delta(\mathcal{H}_0)}{\lambda_0(\mathcal{H}_0)} \geq \frac{1}{n \sqrt{\lambda_{\max}(A)}}$$

"(1st order) Quantum optimization"

? general convex f



This work:

(assume $x^* = 0$, $\nabla f(0) = 0$, $\nabla^2 f(0) = A$)

NON-ASYMPTOTIC

SEMICLASSICAL ANALYSIS

$\gamma \gg 1$

$$H^\gamma = -\frac{1}{2}\Delta + \gamma^2 f$$

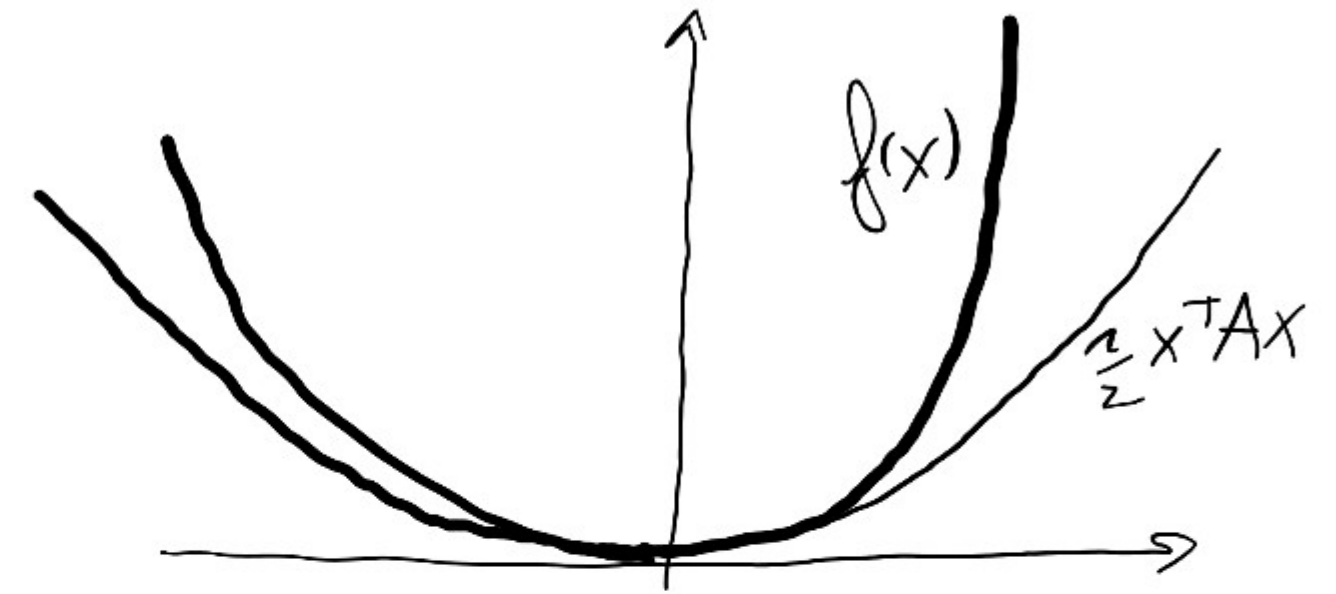
$$\approx -\frac{1}{2}\Delta + \gamma^2 \frac{1}{2} x^T A x = H_0^\gamma$$

for low energies:

$$\lambda_0(H^\gamma) \approx \lambda_0(H_0^\gamma), \quad \lambda_1(H^\gamma) \approx \lambda_1(H_0^\gamma)$$

"(1st order) Quantum optimization"

? general convex f



This work:

(assume $x^* = 0$, $\nabla f(0) = 0$, $\nabla^2 f(0) = A$)

NON-ASYMPTOTIC

SEMICLASSICAL ANALYSIS

"uniform" over family of fcts.

Thm.: If f self-concordant and $\gamma \geq \text{poly}(n, \|A\|, \|A^{-1}\|)$,
then $\lambda_0(H) \approx \gamma \frac{1}{2} \text{Tr}[\sqrt{A}]$ and $s(H) \approx \gamma \sqrt{\|A^{-1}\|}$.

"2nd order quantum optimization"

recoll: gradient descent: $\nabla \leadsto$ Newton step: $H^{-1}\nabla$

! $H^{-1}\nabla$ is "Riemannian gradient"

wrt Hessian manifold $M_g = (\mathbb{R}^n, g = H)$

$$\text{e.g., } d\text{vol}_g(x) = \sqrt{\det H(x)} dx$$

$\leadsto$ we define "2nd order Schrödinger operator"

$$\boxed{H = -\frac{1}{2}L + f}$$

= Schrödinger operator on Hessian manifold

$\hookrightarrow L =$ Laplace-Beltrami operator on M_g

"2nd order quantum optimization"

↳ Schrödinger operator $\mathcal{H} = -\frac{1}{2}L + f$ on Hessian manifold M_g

* key example: quadratic $f(x) = \frac{1}{2}x^T A x \leadsto$ "flat" manifold $M_g = (\mathbb{R}^n, g=A)$

$$\mathcal{H}_0 = -\frac{1}{2}L + f = -\frac{1}{2}\nabla^T A^{-1}\nabla + \frac{1}{2}x^T A x$$

$$= -\frac{1}{2}\Delta_y + \frac{1}{2}y^T y \quad \text{for } y = \sqrt{A}x$$

$$\leadsto \lambda_0(\mathcal{H}_0) = \frac{1}{2}n, \quad \delta(\mathcal{H}_0) = 1 \quad ! \text{ ind. of cond. number}$$

* our work:

Thm.: If f self-concordant and $\gamma \geq \text{poly}(n)$,

then $\lambda_0(\mathcal{H}^\gamma) \approx \gamma^{\frac{1}{2}}n$ and $\delta(\mathcal{H}^\gamma) \approx \gamma$.

$$\mathcal{H}^\gamma = -\frac{1}{2}L + \gamma^2 f \leftarrow$$

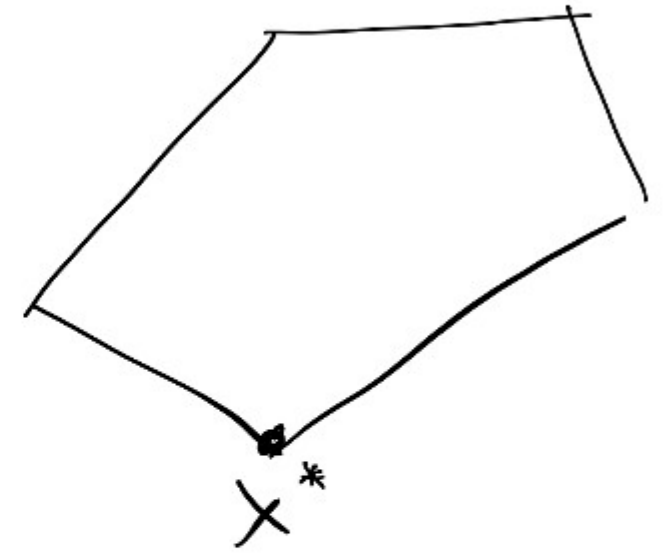
Interior point methods

convex subset $D \subseteq \mathbb{R}^n$

constrained convex optimization:

↳ gradient descent fails!

$$\min_{x \in D} f(x)$$

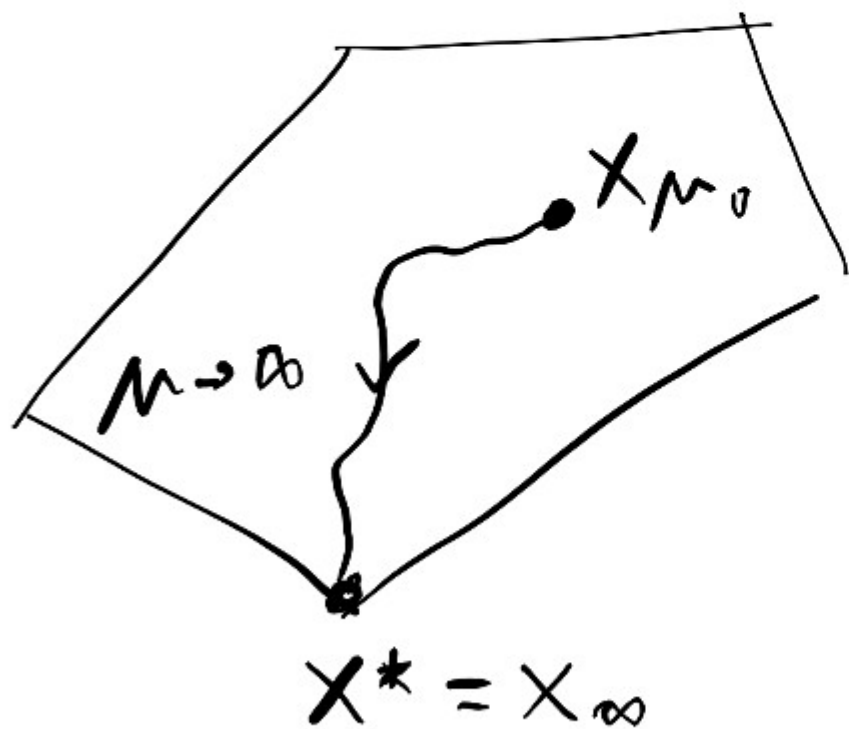


→ interior point methods (IPMs):

1. define barrier function h s.t. $h(x) \rightarrow \infty$ as $x \rightarrow \partial D$

2. solve "unconstrained" problem $\min_x f_\mu(x) = \min_x (\mu f(x) + h(x))$

for increasing μ :



from $x_{\mu_0} = \operatorname{argmin} f_{\mu_0}(x)$,
repeat $\begin{cases} \text{(a) } \mu' = (1+\nu)\mu \\ \text{(b) from } x_\mu, \text{ solve for } x_{\mu'} \end{cases}$

= "path-following"

— using Newton steps

Quantum interior point method

"self-concordant Schrödinger operator"



$$H_\mu := -\frac{1}{2}L + \gamma^2(\mu f + h) \rightarrow \text{g.s. } |\psi_\mu\rangle, \text{ gap } \delta_\mu$$

Assuming (I) initial state $|\tilde{\psi}_{m_0}\rangle \approx |\psi_{m_0}\rangle$

(h) estimates of g.s. energies $\lambda_0(H_\mu)$

(H) (controlled) Hamiltonian simulation e^{iH_μ}

Thm.: From $|\tilde{\psi}_{m_0}\rangle$, prepare $|\psi_\mu\rangle$ concentrated around minimizer x^* while evolving H_μ for time $\text{poly}(n, \log(1/\epsilon))$.

! compared to $\text{poly}(n, \max \kappa(A_\mu), \log(1/\epsilon))$ with 1st order Schrödinger operator

Key ideas

- quantum annealing: $|\psi_0\rangle \rightsquigarrow |\psi_1\rangle \rightsquigarrow \dots \rightsquigarrow |\psi_T\rangle$
 $|\langle\psi_j|\psi_{j+1}\rangle| \in \Omega(1)$ $\nearrow \geq 1/\text{poly}(n)$
- semiclassical analysis: triples $(H_\mu, |\psi_\mu\rangle, \delta_\mu)$
 $\searrow \approx \text{Gaussian}$
- reflection around unbounded operator:

! QPE on e^{iH_μ} fails because of "aliasing"

→ new trick based on "Hubbard-Stratonovich transform":

$$(\mathbb{1} \otimes \langle g |) e^{i t H_\mu \hat{x}} |\varphi\rangle |g\rangle = e^{-\frac{t^2}{2} H_\mu^2} |\varphi\rangle \approx |\psi_\mu \rangle \langle \psi_\mu | \varphi\rangle$$

$\nwarrow$ Gaussian state $\nwarrow$ imaginary-time evolution

Summary:

semiclassical analysis of "second-order" Schrödinger operator,
↕
"continuous-variable" quantum IPM
└ independent of cond. number

Open questions:

- sp. gap bounds using global properties
- discretization: Hamiltonian simulation in curved spaces?